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# **DISCRETE ADJOINT SOLVER FOR DESIGN IN AERONAUTICS**

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## OUTLINE

**1- DISCRETE ADJOINT SOLVER**

**2- RECENT ADVANCES**

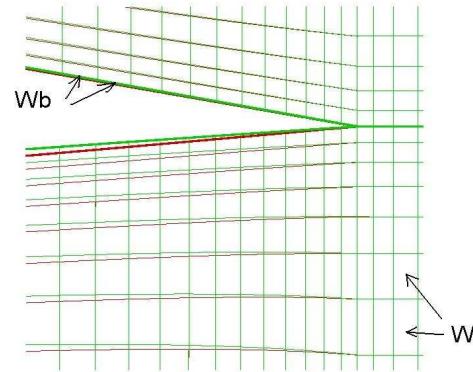
**3- APPLICATIONS**

## Discrete adjoint solver Layout

- Local optimization of aeronautical shapes
- Minimization of an aerodynamic function  $\mathcal{J}$  (drag ...) under the constraints  $\mathcal{G}_k \leq 0$  ( $k \in [1, n_c]$ ) (lift, pitching moment ...).
- Parametric optimization (vector of parameters  $\alpha \in R^p$ ):  $\mathcal{J}(\alpha), \mathcal{G}_k(\alpha)$
- Gradients  $d\mathcal{J}(\alpha)/d\alpha_i$  and  $d\mathcal{G}_{\phi(k)}(\alpha)/d\alpha_i$  ( $i \in [1, n_s]$ ) used in the optimization process.

## Discrete adjoint solver Equations (1/4)

- A finite-volume, cell centered code: elsA <sup>a</sup> (developped at the Onera)
- $X$  is the mesh and  $W$  the aerodynamic conservative field
- $R(W, X)$  is the discrete residual
- $J$  and  $g_k$  are defined such as  $\mathcal{J}(\alpha) = J(W(\alpha), X(\alpha), W_b(X(\alpha), W(\alpha)))$  and  $\mathcal{G}_k(\alpha) = g_k(W(\alpha), X(\alpha), W_b(X(\alpha), W(\alpha)))$



<sup>a</sup> Cambier L, Veuillot J. Status of the elsA software for flow simulation and multi-disciplinary applications. 46th AIAA Aerospace Sciences Meeting and Exhibit, Reno, Nevada, number AIAA Paper 2008-0664, 2008.

## Discrete adjoint solver Equations (2/4)<sup>a</sup>

- Discrete equations for mechanics (set of  $n_a$  non-linear equations )

$$R(W(\alpha), X(\alpha)) = 0$$

- Differentiation with respect to  $\alpha_i$   $i \in [1, n_s]$ . Derivation of  $n_s$  linear system of size  $n_a$

$$\frac{\partial R}{\partial W} \frac{dW}{d\alpha_i} = - \left( \frac{\partial R}{\partial X} \frac{dX}{d\alpha_i} \right)$$

- Calculation of derivatives

$$\nabla_{\alpha} \mathcal{J}(\alpha) = \frac{\partial J}{\partial X} \frac{dX}{d\alpha} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial X} \frac{dX}{d\alpha} + \left( \frac{\partial J}{\partial W} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial W} \right) \frac{dW}{d\alpha}$$

$$\nabla_{\alpha} \mathcal{G}_k(\alpha) = \frac{\partial g_k}{\partial X} \frac{dX}{d\alpha} + \frac{\partial g_k}{\partial W_b} \frac{\partial W_b}{\partial X} \frac{dX}{d\alpha} + \left( \frac{\partial g_k}{\partial W} + \frac{\partial g_k}{\partial W_b} \frac{\partial W_b}{\partial W} \right) \frac{dW}{d\alpha}$$

<sup>a</sup>Peter J, Dwight R. Numerical sensitivity analysis for aerodynamic optimization: A survey of approaches. Comp. Fluids 2010;39:373-391

## Discrete adjoint solver Equations (3/4)

- Following equalities hold  $\forall \lambda \in \mathbf{R}^{n_a}$

$$\lambda^T \frac{\partial R}{\partial W} \frac{dW}{d\alpha_i} + \lambda^T \left( \frac{\partial R}{\partial X} \frac{dX}{d\alpha_i} \right) = 0$$

$$\begin{aligned} \frac{d\mathcal{J}(\alpha)}{d\alpha_i} &= \frac{\partial J}{\partial X} \frac{dX}{d\alpha_i} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial X} \frac{dX}{d\alpha_i} + \left( \frac{\partial J}{\partial W} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial W} \right) \frac{dW}{d\alpha_i} + \\ &\quad \lambda^T \frac{\partial R}{\partial W} \frac{dW}{d\alpha_i} + \lambda^T \left( \frac{\partial R}{\partial X} \frac{dX}{d\alpha_i} \right) \end{aligned}$$

$$\begin{aligned} \frac{d\mathcal{J}(\alpha)}{d\alpha_i} &= \left( \frac{\partial J}{\partial W} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial W} + \lambda^T \frac{\partial R}{\partial W} \right) \frac{dW}{d\alpha_i} \\ &\quad \frac{\partial J}{\partial X} \frac{dX}{d\alpha_i} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial X} \frac{dX}{d\alpha_i} + \lambda^T \left( \frac{\partial R}{\partial X} \frac{dX}{d\alpha_i} \right) \end{aligned}$$

## Discrete adjoint solver Equations (4/4)

- Vector  $\lambda$  defined in order to cancel the factor of the flow sensitivity  $\frac{dW}{d\alpha_i}$

$$\frac{\partial J}{\partial W} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial W} + \lambda^T \frac{\partial R}{\partial W} = 0 \iff \left( \frac{\partial R}{\partial W} \right)^T \lambda = - \left( \frac{\partial J}{\partial W} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial W} \right)^T$$

- Calculation of derivatives

$$\forall i \in [1, nf] \quad \frac{d\mathcal{J}(\alpha)}{d\alpha_i} = \frac{\partial J}{\partial X} \frac{dX}{d\alpha_i} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial X} \frac{dX}{d\alpha_i} + \lambda^T \left( \frac{\partial R}{\partial X} \frac{dX}{d\alpha_i} \right)$$

$$\nabla_\alpha \mathcal{J}(\alpha) = \frac{\partial J}{\partial X} \frac{dX}{d\alpha} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial X} \frac{dX}{d\alpha} + \lambda^T \left( \frac{\partial R}{\partial X} \frac{dX}{d\alpha} \right)$$

- Method with  $(n_c^* + 1)$  (number of functions) and not  $n_s$  linear systems to solve

## Discrete adjoint solver $dJ/dX$ mode<sup>a</sup> (1/3)

- With the classical gradient methods, the storage of  $dX/d\alpha_i$  requires large memory resources  $\Rightarrow dJ/dX$  mode.
- $dJ/dX$  mode: calculate the derivatives of the discrete equations  $R$  w.r.t. the mesh  $X$ , as well as  $dJ/dX$  and  $dG_k/dX$

$$\nabla_{\alpha} \mathcal{J}(\alpha) = \left( \frac{\partial J}{\partial X} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial X} + \lambda^T \frac{\partial R}{\partial X} \right) \frac{dX}{d\alpha}$$

$$\boxed{\nabla_X \mathcal{J}(X) = \frac{\partial J}{\partial X} + \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial X} + \lambda^T \frac{\partial R}{\partial X}}$$

to be calculated and stored.

<sup>a</sup> Trontin P, Peter J, Nguyen-Dinh M. Goal oriented mesh adaptation using total derivative of aerodynamic functions with respect to mesh coordinates. 49th AIAA Aerospace Sciences Meeting, 4-7 January 2011. Accepted.

## Discrete adjoint solver dJ/dX mode (2/3)

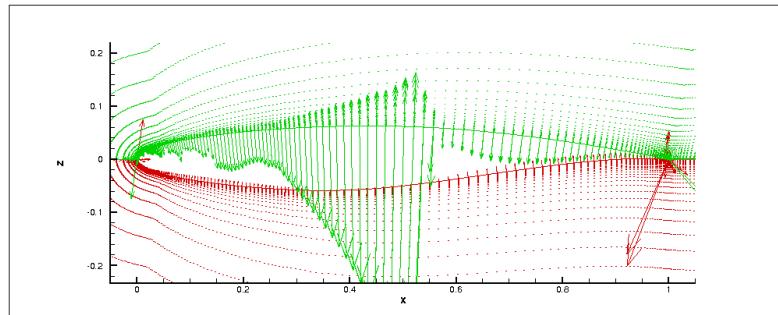
- Final product with  $dX/d\alpha_i$  carried out by another computer with higher memory capacities and possibly lower computational speed.
- Surface mesh  $S(\alpha)$
- Volume mesh  $X$  is an explicit function of  $S$  and  $\tilde{J}(S) = J(X(S))$  (resp.  $\tilde{G}_k(S) = G_k(X(S))$ )

$$\frac{d\tilde{J}}{dS} = \frac{dJ}{dX} \frac{dX}{dS} \quad \frac{d\tilde{G}_k}{dS} = \frac{dG_k}{dX} \frac{dX}{dS}$$

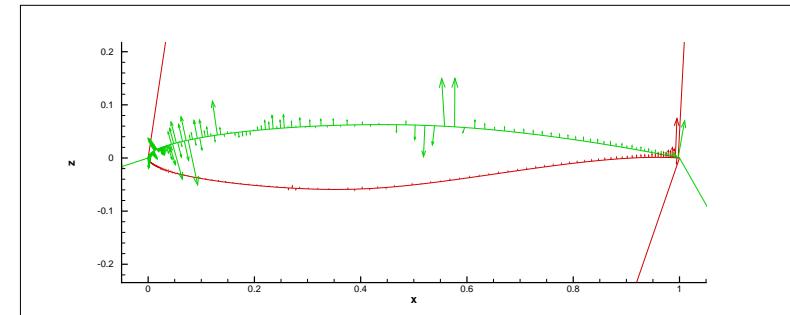
- Shape optimization independant of any set of the design parameters
- A cheaper computation of the sensitivities (multiplying last equation by  $(dS/d\alpha_i)$ )

## Discrete adjoint solver $dJ/dX$ mode (3/3)

- Airfoil RAE2822
- $M_\infty = 0.734, Re_\infty = 6.5 \cdot 10^6$
- AoA =  $2.79^\circ$



$$dC_D/dX$$



$$dC_D/dS$$

# Discrete adjoint solver

## Available Numerical methods in elsA

- Differentiated fluxes

- △ Mean flow :
  - ▷ Roe flux : second order using the MUSCL formula with Van Albada limiting function
  - ▷ Centered flux plus scalar artificial dissipation (Jameson et al., 1981)
  - ▷ Viscous flux with cell-centered gradients corrected at interfaces
- △ Turbulence quantities transport equations
  - ▷ Roe order one flux
  - ▷ Viscous flux
  - ▷ Source term of the differentiated models (cell centered evaluation of gradient terms)<sup>a</sup>
  - ▷ Models :  $k - \omega$  Wilcox,  $k - \epsilon$  Launder-Sharma, Spalart Allmaras

- Resolution of the adjoint equation

- ▷ Newton-type or relaxation algorithm<sup>b</sup>
- ▷ Each Newton step solved by a LU-RELAX or LU-SSOR technique
- ▷ Joins: matching, line matching, plane matching

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<sup>a</sup> Renac F, Pham CT, Peter J. Sensitivity analysis for the RANS equations coupled with a linearized turbulence model. AIAA Paper Series, Paper 2007-3948, 2007

<sup>b</sup> Peter J, Drullion F. Large stencil viscous flux linearization for the simulation of 3D turbulent flows with backward Euler schemes. Comp. Fluids, 36:1007-1027, 2007

# OUTLINE

## 1- DISCRETE ADJOINT SOLVER

## 2- RECENT ADVANCES

- ▷ The full adjoint linearization of the Spalart Allmaras turbulence model
- ▷ The RPM method
- ▷ The chimera method

## 3- APPLICATIONS

# The full linearization of the Spalart Allmaras turbulent model(1/3)

$$|V| \frac{\partial \mathbf{W}}{\partial t} + \oint_{\partial V} (\mathbf{F}_c + \mathbf{F}_v) \cdot \mathbf{n} dS - \int_V \mathbf{S} dV = 0$$

- **conservative variables:**

$$\mathbf{W} = (\rho, \rho \mathbf{U}^t, \rho E, \rho k, \rho \tilde{\nu})^t$$

- **convective and viscous fluxes:**

$$\mathbf{F}_c = \begin{pmatrix} \rho \mathbf{U}^t \\ \rho \mathbf{U} \mathbf{U}^t + p \mathbf{I} \\ (\rho E + p) \mathbf{U}^t \\ \rho \tilde{\nu} \mathbf{U}^t \end{pmatrix}, \quad \mathbf{F}_v = \begin{pmatrix} 0 \\ -\boldsymbol{\tau} - \boldsymbol{\tau}_r \\ -(\boldsymbol{\tau} + \boldsymbol{\tau}_r) \mathbf{U} + \mathbf{q}^t + \mathbf{q}_t^t \\ -\frac{1}{\sigma} (\mu + \rho \tilde{\nu}) \nabla \tilde{\nu}^t \end{pmatrix}$$

- **constraints and heat flux:**

$$\begin{aligned} \boldsymbol{\tau} + \boldsymbol{\tau}_r &= (\mu + \mu_t) \left( -\frac{2}{3} (\nabla \cdot \mathbf{U}) \mathbf{I} + \nabla \mathbf{U} + \nabla \mathbf{U}^t \right) \\ \mathbf{q} + \mathbf{q}_t &= -\left( \frac{C_p \mu}{Pr} + \frac{C_p \mu_t}{Pr_t} \right) \nabla T \end{aligned}$$

## The full linearization of the Spalart Allmaras turbulent model(2/3)

- turbulent viscosity coefficient:

$$\mu_t = \mu \chi f_{v1} \quad \text{with} \quad f_{v1} = \frac{\chi^3}{\chi^3 + C_{v1}^3} \quad \text{and} \quad \chi = \frac{\rho \tilde{\nu}}{\mu}$$

- source terms:

$$\mathbf{S} = (0, 0, 0, Prod + Cross + Dest)^t$$

– production :

$$\begin{aligned} Prod &= C_{b1} \tilde{S} \rho \tilde{\nu} \\ \tilde{S} &= \bar{\omega} + \frac{|\rho \tilde{\nu}|}{\rho} \frac{f_{v2}}{\kappa^2 \eta^2} \\ f_{v2} &= 1 - \frac{\chi}{1 + \chi f_{v1}} \end{aligned}$$

– cross term (diffusion) :

$$\begin{aligned} Cross &= \frac{C_{b2}}{\sigma} \rho \nabla \tilde{\nu} \cdot \nabla \tilde{\nu} \\ Cross &= \min(Cross, 20 \rho \tilde{\nu} \max(0, C_{b1} \tilde{S})) \end{aligned}$$

## The full linearization of the Spalart Allmaras turbulent model(3/3)

- source terms:

- destruction :

$$\begin{aligned}
 Dest &= -C_{w1} f_w \frac{|\rho\tilde{\nu}|}{\rho} \frac{\rho\tilde{\nu}}{\eta^2} \\
 f_w &= g \left( \frac{1 + C_{w3}^6}{g^6 + C_{w3}^6} \right)^{1/6} \\
 g &= \tilde{r} + C_{w2}(\tilde{r}^6 - \tilde{r}) \\
 \tilde{r} &= \max \left[ 0, \min \left( 10, \frac{1}{\tilde{S}} \frac{|\rho\tilde{\nu}|}{\rho} \frac{1}{\kappa^2 \eta^2} \right) \right]
 \end{aligned}$$

- model constants:

coefficient	value	coefficient	value
$C_{b1}$	0, 1355	$C_{w1}$	$\frac{C_{b1}}{\kappa^2} + \frac{1+C_{b2}}{\sigma}$
$C_{b2}$	0, 622	$C_{w2}$	0, 3
$\sigma$	$\frac{2}{3}$	$C_{w3}$	2
$\kappa$	0, 41	$C_{v1}$	7, 1

## Results on the airfoil RAE2822

- Reynolds number:  $Re = \frac{\sqrt{\rho_\infty p_\infty} c}{\mu(T_\infty)} = 6.5 \times 10^6$
- Mach number:  $M = \frac{U_\infty}{c_\infty} = 0,73$
- AoA:  $\alpha = 2.79$  deg.
- spatial discretization : 2 blocks of  $73 \times 229$  knots
- rotation variation (for finite differences) :  $\delta\alpha = 10^{-2}$  deg.
- design parameter: angle of rotation around ( $Oy$ )



Figure 1: Airfoil RAE2822

## Adjoint gradient computation on the airfoil RAE2822

function	DF	frozen turbulence assumption	full linearization
$dC_{L,p}/d\alpha$	1,703e-01	1,105e-01 (35%)	1,687e-01 (0, 9%)
$dC_{D,p}/d\alpha$	5,688e-03	4,741e-03 (17%)	5,703e-03 (0, 3%)
$dC_{D,w}/d\alpha$	2,272e-03	1,658e-03 (27%)	2,251e-03 (0, 9%)
$dC_{D,ff}/d\alpha$	4,092e-03	3,606e-03 (12%)	4,069e-03 (0, 5%)
$dC_{D,vp}/d\alpha$	1,820e-03	1,948e-03 (7%)	1,818e-03 (0, 1%)

# OUTLINE

## 1- DISCRETE ADJOINT SOLVER

## 2- RECENT ADVANCES

- ▷ The full linearization of the Spalart Allmaras turbulent model
- ▷ The RPM method
- ▷ The chimera method

## 3- APPLICATIONS

## Backward Euler scheme

- The adjoint equation to be solved:

$$\left( \frac{\partial R}{\partial W} \right)^{(ACC)T} \lambda = - \left( \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial W} + \frac{\partial J}{\partial W} \right)^T$$

- For solving the linear system:

$$\mathbf{Ax} = \mathbf{b}$$

with  $\mathbf{A} = \left( \frac{\partial R}{\partial W} \right)^{(ACC)T}$ ,  $\mathbf{x} = \lambda$  and  $\mathbf{b} = - \left( \frac{\partial J}{\partial W_b} \frac{\partial W_b}{\partial W} + \frac{\partial J}{\partial W} \right)^T$

- we use the iterative algorithm:

$$\mathbf{x}^{(l+1)} = \mathbf{F}(\mathbf{x}^{(l)}) = \Phi \mathbf{x}^{(l)} + \mathbf{M}^{-1} \mathbf{b}$$

– iteration matrix:

$$\Phi = \mathbf{I} - \mathbf{M}^{-1} \mathbf{A}$$

– implicit matrix:

$$\mathbf{M} = \frac{1}{\Delta t} \mathbf{I} + \frac{\partial \mathbf{R}}{\partial \mathbf{W}}^{(APP)}$$

# Recursive Projection Method (RPM)<sup>a</sup>

## Stability of the iterative scheme

- Theorem: convergence of the iterative algorithm

Let  $\mathbf{M}, \mathbf{A} \in \mathbb{C}^{N \times N}$  nonsingular, then

$$\forall \mathbf{x}^{(0)}, \quad \lim_{l \rightarrow \infty} \mathbf{x}^{(l)} = \mathbf{A}^{-1}\mathbf{b} \quad \Leftrightarrow \quad \rho(\Phi) < 1$$

- Divergent eigenspace definition

- divergent eigenmodes  $(\lambda_i, \mathbf{e}_i)$ :  $|\lambda_1| \geq \dots \geq |\lambda_m| \geq 1$
- divergent eigenspace:

$$\mathbb{P} = \text{span}\{\mathbf{e}_1, \dots, \mathbf{e}_m\} ; \quad \mathbb{Q} = \mathbb{P}^\perp$$

- orthogonal projectors:

$$\forall \mathbf{x} \in \mathbb{R}^N, \quad \mathbf{x} = \mathbf{P}\mathbf{x} + \mathbf{Q}\mathbf{x} = \mathbf{x}_p + \mathbf{x}_q$$

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<sup>a</sup>Renac F. Aerodynamic sensitivity analysis of the RANS equations via the recursive projection method. AIAA Paper Series, Paper 2010-4364, 2010

# Recursive Projection Method (RPM) Stabilization procedure

Let  $\mathbf{V} \in \mathbb{R}^{N \times m}$  be an orthonormal basis for  $\mathbb{P}$ :

$$\mathbf{P} = \mathbf{V}\mathbf{V}^T, \quad \mathbf{Q} = \mathbf{I} - \mathbf{V}\mathbf{V}^T, \quad \mathbf{V}^T\mathbf{V} = \mathbf{I}$$

- RPM algorithm

$$\begin{aligned}\mathbf{x}_q^{(l+1)} &= (\mathbf{I} - \mathbf{P})\mathbf{F}(\mathbf{x}_p^{(l)} + \mathbf{x}_q^{(l)}) \\ \mathbf{x}_p^{(l+1)} &= \mathbf{x}_p^{(l)} + (\mathbf{I} - \mathbf{P}\Phi\mathbf{P})^{-1}(\mathbf{P}\mathbf{F}(\mathbf{x}_p^{(l)} + \mathbf{x}_q^{(l)}) - \mathbf{x}_p^{(l)}) \\ \mathbf{x}^{(l+1)} &= \mathbf{x}_p^{(l+1)} + \mathbf{x}_q^{(l+1)}\end{aligned}$$

- Theorem: unconditional convergence of the RPM iteration

Let  $\mathbf{M}, \mathbf{A} \in \mathbb{C}^{N \times N}$  nonsingular and such that  $\mathbf{P}\Phi\mathbf{Q} = 0$ , then

$$\forall \mathbf{x}^{(0)}, \quad \lim_{l \rightarrow \infty} \mathbf{x}_p^{(l)} + \mathbf{x}_q^{(l)} = \mathbf{A}^{-1}\mathbf{b}$$

# Recursive Projection Method (RPM)

## Krylov acceptance criterion

**Goal:** construction of the divergent basis  $\mathbf{V}$

**Method:**

- let  $k \in \mathbb{N}$
- define the Krylov subspace generated by  $\mathbf{v}_1 = \Delta \mathbf{x}_q^{(l-k+1)}$  and  $\hat{\mathbf{A}} = \mathbf{Q}\Phi\mathbf{Q}$ :

$$\mathbb{K}_k = \text{span}\{\mathbf{v}_1, \hat{\mathbf{A}}\mathbf{v}_1, \dots, \hat{\mathbf{A}}^{k-1}\mathbf{v}_1\}$$

$$\text{where } \Delta \mathbf{x}_q^{(j)} = \hat{\mathbf{A}}^j \mathbf{x}_q^{(0)} = \mathbf{x}_q^{(j+1)} - \mathbf{x}_q^{(j)}.$$

- set

$$\mathbf{K}_k = (\mathbf{v}_1, \hat{\mathbf{A}}\mathbf{v}_1, \dots, \hat{\mathbf{A}}^{k-1}\mathbf{v}_1)$$

- and compute the QR factorization

$$\mathbf{K}_k = \mathbf{Q}_k \mathbf{R}_k$$

# Recursive Projection Method (RPM) Application to Onera M6 wing (1/2)<sup>a</sup>

- Test-case : M6 wing.  $Re = 1.46 \cdot 10^7$ ,  $M_\infty = 0.84$ ,  $AoA = 0^\circ$
- RANS equations + Spalart-Allmaras
- Structured mesh with 10 blocks and 620 042 nodes
- Design parameter: angle of rotation around  $(O, y)$

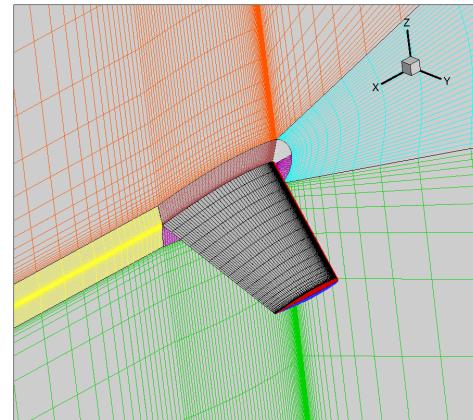
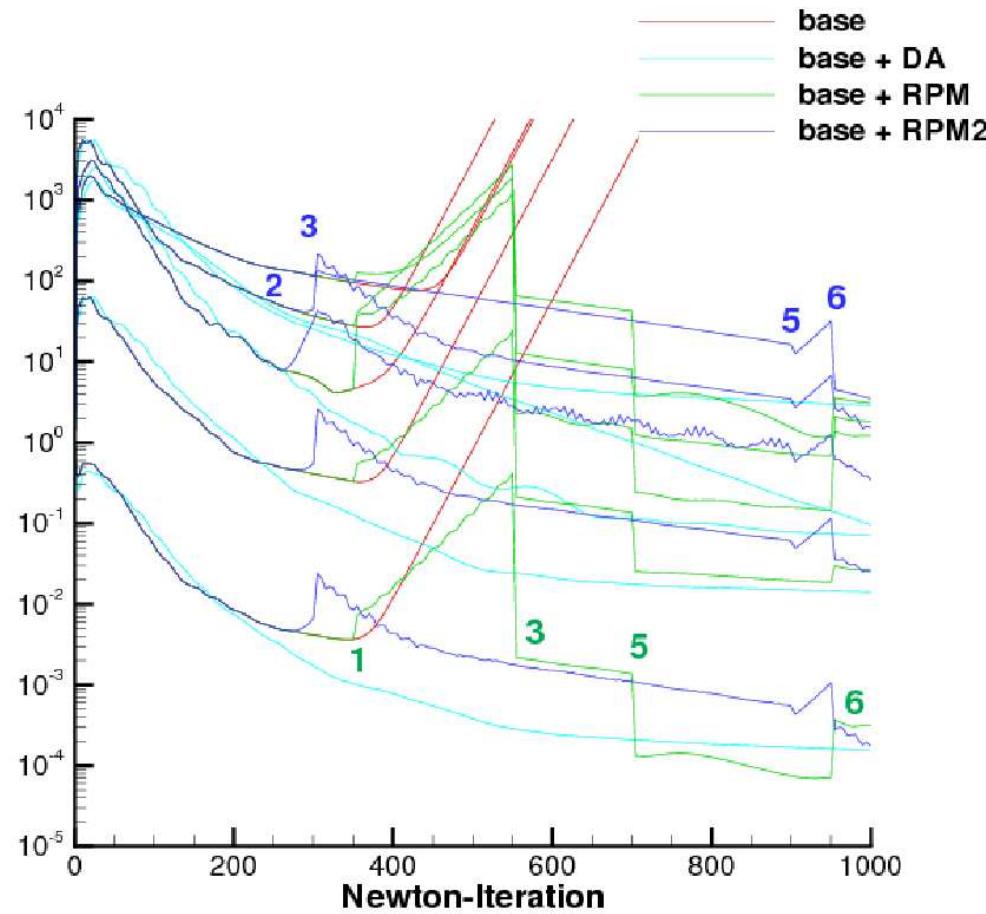


Figure 2: M6 wing

<sup>a</sup> Renac F. Aerodynamic sensitivity analysis of the RANS equations via the recursive projection method. AIAA Paper Series, Paper 2010-4364, 2010

## Recursive Projection Method (RPM) Application to Onera M6 wing (2/2)



# OUTLINE

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## 2- RECENT ADVANCES

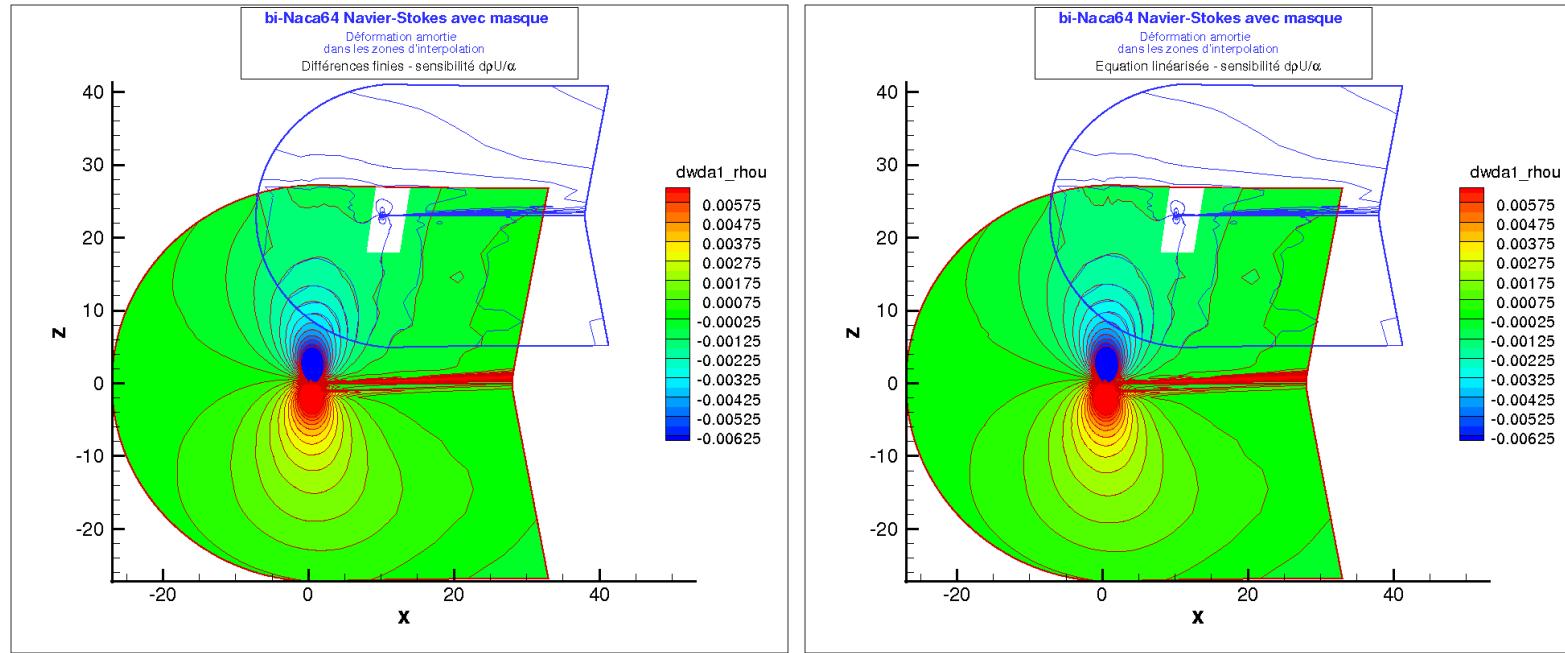
- ▷ The full linearization of the Spalart Allmaras turbulent model
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## 3- APPLICATIONS

## The chimera method (1/2)

- Bi-NACA64A212 with mask
- 2D RANS equations.  $M_\infty = 0.71$ ,  $AoA = 0.25^\circ$ ,  $Re = 2 \cdot 10^6$
- 2 blocks:  $257 \times 63$  and  $257 \times 65$
- Design parameter: camber of the trailing edge of the upper profile
- Limit: blanking and interpolation zones not deformed by the design parameter

## The chimera method (2/2)



	finite difference	adjoint equation	relative error
$dCLp/d\alpha$	-0.5099	-0.5220	2.37 %
$dCDp/d\alpha$	$-1.411 \cdot 10^{-2}$	$-1.494 \cdot 10^{-2}$	5.88 %

## OUTLINE

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**2- RECENT ADVANCES**

**3- APPLICATIONS**

# Aircraft <sup>a</sup> <sup>b</sup>

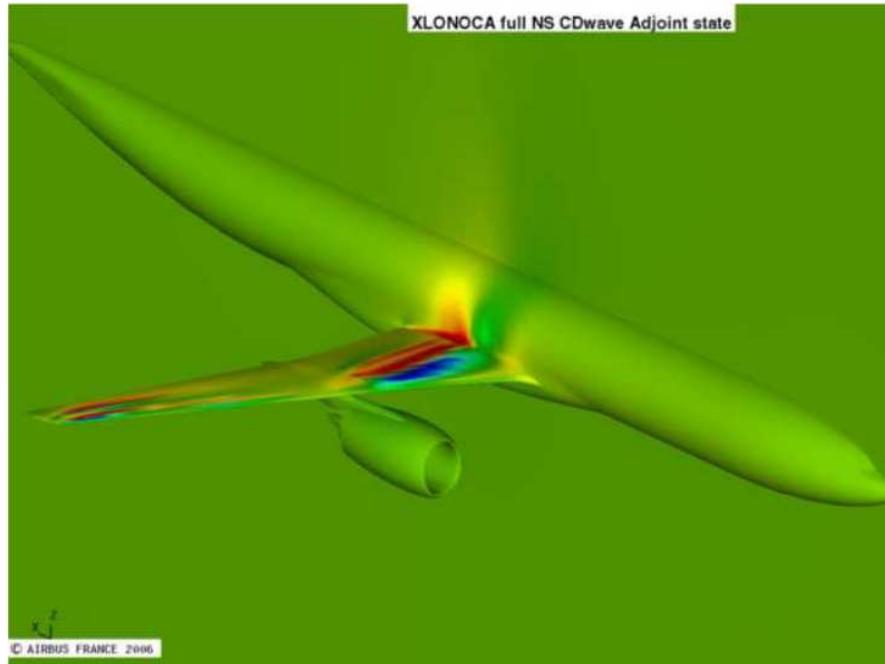


Figure 3: Adjoint vector component for  $CDw$

<sup>a</sup> Salah El Din I, Carrier G, Mouton S. Discrete adjoint method in elsA (Part 2): Application to aerodynamic design optimization. Proceedings of the 7th ONERA-DLR Aerospace Symposium (ODAS), Toulouse, 2006

<sup>b</sup> Carrier G. Single and multi-point aerodynamic optimizations of a supersonic transport aircraft wing using optimization strategies involving adjoint method and genetic algorithm. Proceedings of ERCOFAC Workshop "Design optimization: methods and applications", Las Palmas, 2008

# Turbomachinery

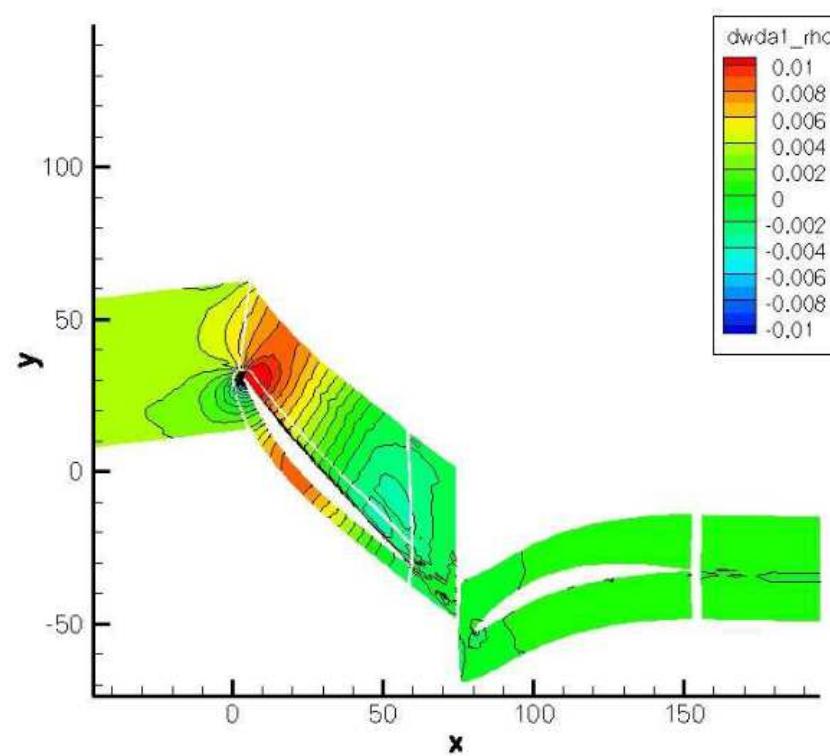


Figure 4: CME2 3D, rotor stator with clearance.  $d\rho/d\alpha$

## Rotor in hover (1/2)<sup>a</sup>

- Rotor ERATO
- 3D RANS equations.  $M_{tip} = 0.616$ ,  $Re_{tip} = 1.93 \cdot 10^6$
- 1 block: 807 457 nodes
- Design parameters: poles of Bezier curve driving chord, twist and sweep

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<sup>a</sup>Dumont A, Le Pape A, Peter J, Huberson S. Aerodynamic shape optimization of hovering rotors using a discrete adjoint of the RANS equations. Proceedings of 5th American Helicopter Society annual forum, Grapevine, Texas, USA, 2009. To appear in AHS Journal.

## Rotor in hover (2/2)

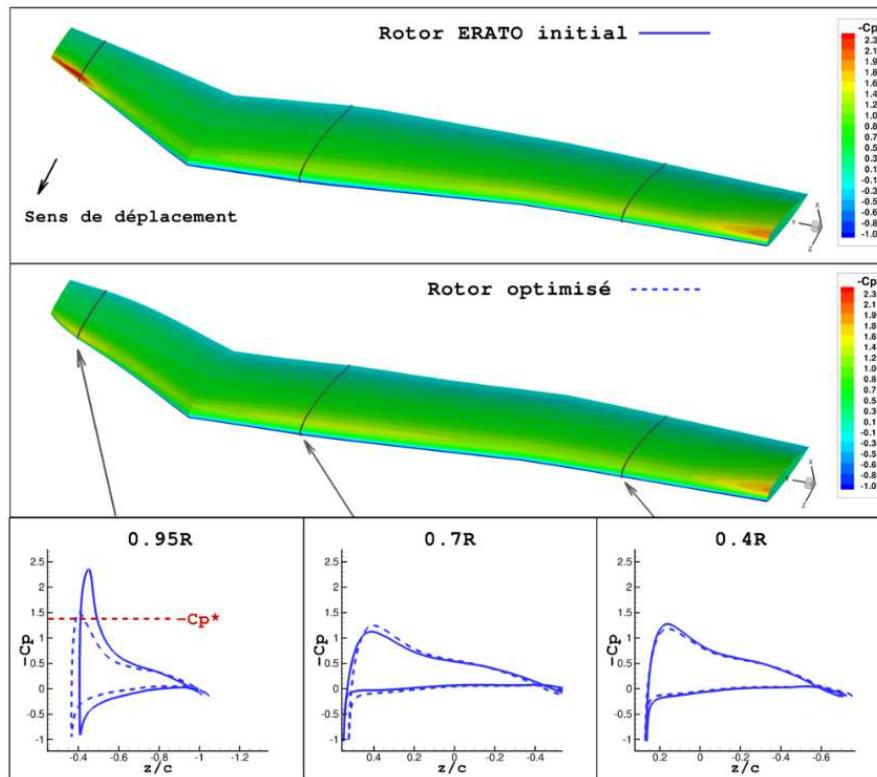


Figure 5: Rotor ERATO

**THANK YOU FOR YOUR ATTENTION**